

$$5. \quad u(x,0) = \begin{cases} 0 & x < 0 \\ 10-x & x \geq 0 \end{cases}$$

This is not applicable.

$$u(x,t) = \int_{-\infty}^{\infty} \Phi(x-y,t) \cdot u(y,0) dy$$

$$= \int_{-\infty}^{\infty} \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-y)^2}{4kt}} u(y,0) dy$$

$$= \int_0^{\infty} \frac{1}{\sqrt{4\pi kt}} e^{-\frac{(x-y)^2}{4kt}} (10-y) dy$$

$$\text{Erf}(x) = \frac{2}{\sqrt{\pi}} \cdot \int_0^x e^{-y^2} dy$$

$$z = x - y$$

Substitute  $x-y$  by  $z$  ("="  $y = x-z$ )

$$= \int_x^{-\infty} \frac{1}{\sqrt{4\pi kt}} \cdot e^{-z^2/4kt} \cdot (10 - x + z) dz$$

$$= \int_{-\infty}^x " \cdot e^{-z^2/4kt} \cdot (10 - x + z) dz$$

$$\int e^{-z^2/4kt} \cdot z dz = \text{const.} \cdot e^{-z^2/4kt}$$

$$(10-x) \cdot \underbrace{\int_{-\infty}^x e^{-\frac{z^2}{kbt}} dz}_{\text{"const" w.r.t. } z} = \underbrace{\int_{-\infty}^0}_{\downarrow \text{explicit number}} + \underbrace{\int_0^x}_{\downarrow \text{Erf}}$$

$$\Rightarrow u(x,t) = \text{const.} \cdot e^{-\frac{x^2}{kbt}} + \boxed{\text{some number}} \cdot (10-x) + \text{Erf}(\quad) \cdot (10-x).$$

$$6. \quad u(x,0) = x^2.$$

$$(a) \quad \text{if } u \text{ satisfies } u_t = k u_{xx}$$

$$\text{using Clairaut's thm, } u_{xxx} = k \cdot \frac{u_{xxxx}}{u_t} \quad || \\ u_{xxx} = \frac{u_{xxxx}}{u_t}.$$

$$v(x,0) = u_{xxx}(x,0) = (x^2)''' = 0.$$

$$(b) \quad \text{Fact: if } v(x,0) = 0, \text{ then } v(x,t) = 0.$$

$$\Rightarrow u_{xxx}(x,t) = 0.$$

$$\Rightarrow u(x,t) = C_2 \cdot x^2 + C_1 \cdot x + C_0.$$

$C_2, C_1, C_0$  are functions of  $t$ .

$$\textcircled{1} \quad u_t = k u_{xx} \quad \begin{matrix} \parallel \\ 0 \end{matrix} \quad \begin{matrix} \nearrow x^2 \\ \parallel \\ 0 \end{matrix} \quad \begin{matrix} \nearrow x \\ \parallel \\ 0 \end{matrix}$$

$$\Rightarrow (\text{LHS}) \quad C_2'(t) \cdot x^2 + C_1'(t) \cdot x + C_0'(t)$$

$$(\text{RHS}) \quad \boxed{k \cdot C_2(t) \cdot 2} \quad \xrightarrow{\text{Constant part}}$$

$$\textcircled{2} \quad u(x, 0) = x^2$$

$$\Rightarrow C_2(0) \cdot x^2 + C_1(0) \cdot x + C_0(0) = x^2$$

$$\begin{matrix} \parallel \\ 1 \end{matrix} \quad \begin{matrix} \parallel \\ 0 \end{matrix} \quad \begin{matrix} \parallel \\ 0 \end{matrix}$$

$\Rightarrow$  Graph info about  $C_2(t), C_1(t), C_0(t)$   
 $\Rightarrow$  can solve for them.

7. Conservation of energy for the wave egn.  
 (Dirichlet condition or Neumann condition.)

Statement : Let  $u$  be a solution for the wave egn.

$$u_{tt} = c^2 \Delta u \quad \text{w/ } u(x, 0) = \phi(x)$$

$$\text{and } u_t(x, 0) = \psi(x)$$

for  $x \in U, t > 0$

Moreover, assume that  $u$  satisfies the D or N  
 boundary condition.  
 homogeneous

lets define the energy function  $E(t)$  as follows:

$$E(t) := \int_U \left( \frac{1}{2} u_t^2 + \frac{c^2}{2} \left( \sum_{i=1}^n u_i^2 \right) \right) dV$$

where  $u_i$  denotes  $\frac{\partial u}{\partial x_i}$ . Then,  $E(t)$  is constant.

pf.  $\frac{d}{dt} E(t) = \int_U \frac{\partial}{\partial t} \left( \frac{1}{2} u_t^2 + \frac{c^2}{2} \left( \sum u_i^2 \right) \right) dV$

$$= \int_U \left( u_t \cdot u_{tt} + c^2 \sum_{i=1}^n (u_i \cdot u_{it}) \right) dV$$

$$c^2 \Delta u = c^2 \sum_{i=1}^n u_{ii} \quad (\text{$u_{ii}$ means } \frac{\partial^2 u}{\partial x_i^2})$$

$$= \int_U c^2 \sum_{i=1}^n (u_t \cdot u_{ii} + u_i \cdot u_{it}) dV$$

$$= \int_U c^2 \cdot \sum_{i=1}^n (u_t \cdot u_{ii}) \downarrow \text{Clairaut} dV$$

$$= c^2 \cdot \int_U \sum_{i=1}^n \frac{\partial(u_t u_{ii})}{\partial x_i} dV \quad n \leftarrow n\text{-dim'l volume factor}$$

Divergence theorem

$$= c^2 \cdot \int_{\partial U} \boxed{(u_t u_1, \dots, u_t u_n) \cdot \vec{n}} dV_{n-1}$$

1) (hom) Dirichlet boundary

$$\Rightarrow u_t(x, t) = 0 \text{ for } x \in \partial U.$$

$$\begin{aligned} & \downarrow \\ &= C^2 \cdot \int_{\partial U} (0, 0, \dots, 0) \cdot \vec{n} dV_{n-1} \\ &= 0. \end{aligned}$$

2) (hom) Neumann boundary.

$$\underline{\text{grad } u_t \cdot \vec{n} = 0}$$

$$\frac{\partial u_t}{\partial v} = 0$$

$\downarrow$   $v$ : normal direction  
to  $\partial U$ .

$$(u_1, \dots, u_n) \cdot v = 0.$$

$$= C^2 \cdot \int_{\partial U} 0 dV_{n-1} = 0.$$

$$\frac{d}{dt} E(t) - \dots = 0 \Rightarrow E(t) \text{ is constant.}$$

8. Suppose  $u_1$  and  $u_2$  are two sol'n's.

Define  $v(x,t) := \underline{u_1(x,t) - u_2(x,t)}$ .

Then,  $\underline{v_t} = k \underline{v_{xx}}$

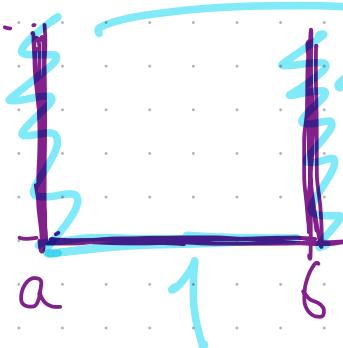
$$\boxed{v(x,0) = u_1(x,0) - u_2(x,0) = 0} \quad 1$$

$$v_t(x,t) =$$

This is a heat egn for a bounded interval  $[a,b]$ .

$\Rightarrow$  Use maximum principle!

$$\underline{v(x,t) \leq 0} \text{ for all } x \in [a,b], t > 0.$$

$t_0$   Dirichlet boundary condition.

but we can also consider

$$u_2 - u_1 = -v$$

$$\Rightarrow -v(x,t) \leq 0.$$

$$\Rightarrow v(x,t) = 0$$

$$\dots \Rightarrow u_1 = u_2 \quad (\text{OR } u_1 - u_2 = 0)$$

$\Rightarrow$  the sol'n is unique

q. Define  $\omega(x,t) = u(x,t) - v(x,t)$ .

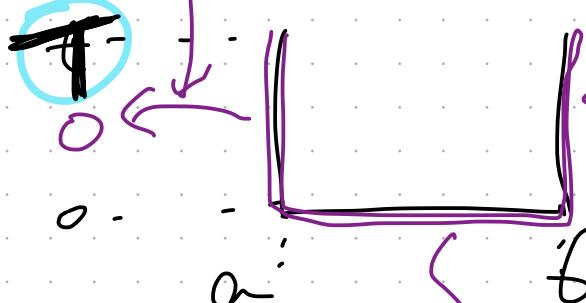
$\omega$  is a soln of the heat eqn.

$$\hookrightarrow \textcircled{1} \quad \omega_t = k \omega_{xx}$$

$$\textcircled{2} \quad \omega(a,t) = \omega(b,t) = 0$$

minimum principle.

$$\Rightarrow \min_{\mathbb{T}} \omega(x,t)$$



is obtained on  $(-\infty, \infty)$

$\Rightarrow$  Tychonoff's solution.

$(\omega(x,t) = u(x,t) - v(x,t))$

$$\Rightarrow \underline{\omega(x,t) \geq 0} \quad \text{for any } t < T$$

but  $T$  is arbitrary  $\Rightarrow \omega(x,t) \geq 0$

$\hookrightarrow t > 0$ .

"homogeneous heat equation"

$\Rightarrow u$  satisfies

$$\textcircled{1} \quad u_t = k u_{xx}$$

$$\textcircled{2} \quad u(x,0) = \phi(x)$$

$$\textcircled{3} \quad u(a,t) = \psi_a(t) \quad u(b,t) = \psi_b(t)$$

This is true  
only for  $[a,b]$ .

$(-\infty, \infty)$

$\Rightarrow$  this will fail.